

# Crystal Plasticity Modeling of Al Alloy under Linear and Non-Linear Loading

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**Abstract.** The crystal plasticity (CP) model is widely used in many applications to link microstructure and mechanical properties. There are varying CP constitutive laws with phenomenological or physical-based formulation to cover a large range of loading conditions. In order to predict the deformation behavior of an Al alloy during the sheet metal forming process with either linear or non-linear strain path, both phenomenological and physical-based CP constitutive laws have been chosen, and the prediction performance of both models is compared. For the linear loading condition, the uniaxial tensile tests are performed on the smooth-dog-bone (SDB) specimens along rolling and transverse directions (RD/TD). The non-linear strain path is achieved by the Marciniak testing followed by uniaxial tension. In the first stage, the Marciniak testing is performed under the stress states of RD-uniaxial, plane strain, and biaxial tension. After being loaded to a certain strain level, mini-SDB specimens are cut along RD and TD from the uniform deformation region and reloaded under RD-uniaxial tension. The digital image correlation (DIC) technique is employed to measure the strain during testing. The electron backscatter diffraction (EBSD) technique is used to characterize the initial microstructure as well as the microstructure evolution of the specimens after the first stage loading in the non-linear strain path. A phenomenological power law and a dislocation-density-based hardening law have been employed in this study. The parameters are calibrated based on the flow curve of the RD uniaxial tension. The model performance is validated by stress-strain response under all the rest loading conditions including the non-linear loading path.

## 1 Introduction

In the last years, aluminum alloys have become widely used in the automotive industry. The main advantages of these materials are a high strength-to-weight ratio combined with good corrosion resistance. Aluminum is increasingly used in structural parts, body panels and various other components, and AA5754 alloy was chosen as the material for this study [1]. It is well known that the mechanical properties of metal materials are determined by the microstructure. Therefore, learning about the quantitative relationship between the microstructure and mechanical properties is very important [2]. Particularly, metal will be subjected to various linear or non-linear loading in the process of forming and application, which will cause different stress states to it and change its mechanical properties [3]. Giving materials different loading and studying the response behavior is of high interest for the sheet metal production and forming industry.

There has been a large number of studies focusing on bridging the microstructure and macroscopic mechanical properties by both experimental and numerical methods. The interested microstructural features [4] normally are the phase fraction, grain morphology, crystal orientation, secondary phase

morphology, etc. While the focused mechanical properties can be various from microscopic level strain partition, macroscopic flow behavior, to fracture and fatigue properties, etc. For example, the effect of texture and dislocation structure on strain hardening anisotropy of aluminum alloy was studied by Lopes, et al [5]. J. W. Yoon, et al. [6] studied finite element simulations of the simple shear test were conducted for 1050-O and 6022-T4 aluminum alloy sheet samples. J. Wang, et al. [7] developed a crystal plasticity finite element model accounting for the microstructural features for simulating the fretting fatigue of AA7075-T651. However, for aluminum alloys that can be considered as a single phase, the challenges that hinder this study are mainly from the following two aspects: i) necessary to synthesize fine microstructure models, accurately characterize the microstructure of materials and accurately control the microstructure variables, and ii) two models are used to conduct efficient and reliable calibration of aluminum alloy under different stress conditions. Lian, et al. [8] developed a calibration strategy for synthetic microstructure and crystal plasticity parameters for producing fine grain structure aluminum alloys, which could systematically and quantitatively analyze the influence of microstructure characteristics.

This study aims to validate these complex pathways using the current formulation of either phenomenological constitutive models or the mechanism-based crystal plasticity models. We are intending to extend the experimental envelope to more general dimensions by combining the change of stress states and the loading angles, which are currently separated into two different research lines. This would create more general conditions that the data generated could be used to further develop, calibrate, and validate both the phenomenological models and mechanism-based models. Further analysis of the microstructure changes in this direction is also helpful to understand the plastic deformation mechanism from the microstructure and dislocation levels. A general research strategy is to decompose a complex path into several separate strain paths to simplify the problem [9]. Considering the mechanical properties of the material, the stress state and the loading direction are considered in the loading process. Generally, the loading design of the two stages is as follows. In the first stage, uniaxial, plane strain, and biaxial pre-stretching can be carried out on the specimen, and then small uniaxial tensile specimens are cut from the specimen after initial loading along different loading directions for the second step of the tensile test.

## 2 Material and Experiments

### 2.1. Material characterization

The material used in this study is a cold-rolled AA5754 H111 Aluminum alloy sheet of 1.5 mm thickness. H111 stands for the tempering type, and it recognizes that the alloy underwent some amount of cold strain hardening after annealing. The electron backscattered diffraction (EBSD) technique has been used to characterize the microstructure of the material, following the strategy used by Liu et al. [10]. The EBSD measurements were performed at 20 kV. The microstructure features like grain size, shape, and texture have been characterized. Samples have been obtained on two planes, i.e. the rolling direction (RD)-transverse direction (TD) plane and the RD-normal direction (ND) plane. In addition, varying investigation areas were chosen through thickness on the RD-ND plane to cover the possible thickness gradient. Grain reconstruction and analysis have been carried out by MATLAB/MTEX toolbox [11]. 15° misorientation threshold was used for the grain reconstruction. The investigated material has an average grain size of 16  $\mu\text{m}$  and mean grain shape aspect ratio of 1:0.681:0.662. The typical fcc rolling texture is observed with a texture index of 1.24. Detailed information microstructure characterization methods and results is referred to Lian et al. [8].

### 2.2. Mechanical testing program

In this study, two categories of loading path experiments were designed, i.e. linear loading and non-linear loading respectively. Linear loading experiment for anisotropy investigation including the tensile test along different loading angles with respect to RD, i.e. 0°, 15°, 30°, 45°, 60°, 75°, and 90°, as shown in Fig. 1 (a). Quasi-static loading is applied with a constant strain rate of  $10^{-4} \text{ s}^{-1}$ . For the non-linear strain paths loading experiment, the material has been firstly pre-strained under three stress states (i.e. RD-uniaxial, plane strain, and biaxial tension) using the hydraulic press and the Marciniak

punch, as shown in Fig. 1 (b) and (c). After being loaded to a certain strain level, mini smooth dog-bone (SDB) specimens were cut along RD and TD from the uniform deformation region of the Marciniak samples and reloaded under uniaxial tension. The dimensions of mini smooth dog-bone (SDB) tensile specimens are given in Fig. 1 (d). To prevent the possible edge crack formation at the edge of the cup before the targeted strain of the planar homogenous area has been reached, the stroke of the press was limited to just below the point of edge fracture. The obtained pre-strains were measured by the GOM Aramis DIC system. For each loading condition, 3 parallel tests were conducted to guarantee repeatability.

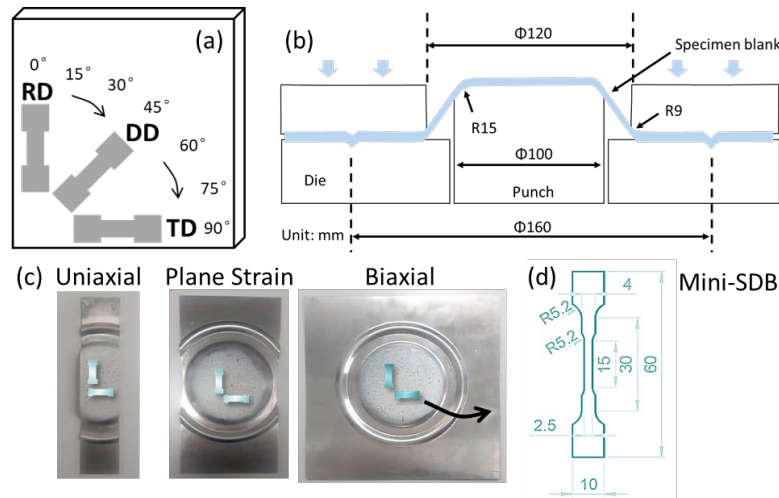


Fig. 1 Experimental procedure of the linear (a) and non-linear strain path: (b) Marciniak punch; (c) three stress states (uniaxial, plane-strain, and biaxial tensions); (d) mini smooth dog-bone tensile test specimens. (unit: mm).

### 3 Numerical Simulations Set Up

#### 3.1 Artificial microstructure model generation

To simulate the microstructure of the chosen material, a representative volume element (RVE) has been employed with the statistically characterized microstructure features in terms of grain size, shape, and texture. This RVE is created by the software DREAM3D and can be adjusted in size and resolution. In our previous study [8], further optimization and filtering are needed to avoid exaggerated deviation of microstructure characteristics. Therefore, the microstructural representative evaluation criterion (MRAC), which can reduce the deviation to 0 [10]. It is defined in the equation below:

$$\Delta = \frac{1}{n} \sum_{i=1}^{i=n} \frac{|X_{i,RVE} - X_{i,ref}|}{X_{i,ref}} \quad (1)$$

Where  $X$  is the characteristic values, and  $i$  represents the type of the microstructure features, and  $n$  is the total number of the characteristic values involved. The  $X_{ref}$  and  $X_{RVE}$  are the characterized values of the corresponding microstructural feature  $X$  from RVE input and output, respectively.

During the RVE generation process, the most representative and effective factors are element size and number, which are corresponding to the final RVE size. The finer element size and a large RVE will have a better representative but they are also will cause a higher computational cost. Considering the balance of RVE representativeness and the computational performance, the final element number is 64,000 elements with 40 elements per edge of the RVE, the element size is 3  $\mu\text{m}$  and the whole RVE size is 120x120x120  $\mu\text{m}^3$ , respectively. The finally optimized RVE model is shown in Fig. 2, which contains 3165 Al grains. The detailed comparisons between optimal RVE and the reference material microstructure are referred to Lian, et al. [8]

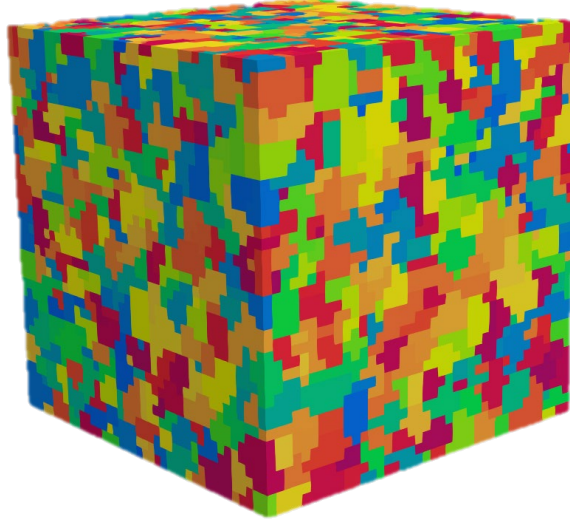


Fig. 2 Grain maps of the optimal RVE.

### 3.2 Phenomenological crystal plasticity model

The crystal plasticity model and the governing equations are based on the dislocation slip mechanism theory [12]. In the phenomenological crystal plasticity model, the slip rate  $\dot{\gamma}^\alpha$  is calculated by the kinetic law of the slip plane system  $\alpha$  (Eq. 2).

$$\dot{\gamma}^\alpha = \dot{\gamma}_0 \left| \frac{\tau^\alpha}{\tau_c^\alpha} \right|^m \text{sgn}(\tau^\alpha) \quad (2)$$

$\dot{\gamma}_0$  is the reference shear rate and  $m$  the rate sensitivity.  $\tau^\alpha$  is the resolved shear stress on the slip system  $\alpha$  and  $\tau_c^\alpha$  the critical resolved shear stress.  $\tau^\alpha$  is defined by Eq. 3.

$$\tau^\alpha = \mathbf{S} \cdot (\mathbf{m}^\alpha \otimes \mathbf{n}^\alpha) \quad (3)$$

$\mathbf{S}$  is the second Piola-Kirchhoff stress in the intermediate configuration. The micromechanical interaction between different slip systems is also given by Eq. 4.

$$\dot{\tau}_c^\alpha = \sum_{\beta=1}^N h_{\alpha\beta} |\dot{\gamma}^\beta| \quad (4)$$

$h_{\alpha\beta}$  is the hardening matrix and is given by Eq. 5.

$$h_{\alpha\beta} = q_{\alpha\beta} \left[ h_0 \left( 1 - \frac{\tau_c^\beta}{\tau_c^s} \right)^a \right] \quad (5)$$

$h_0$ ,  $a$  and  $\tau_c^s$  are the slip hardening parameters.  $q_{\alpha\beta}$  describes the effect of self-hardening ( $\alpha = \beta$ ) and latent hardening ( $\alpha \neq \beta$ ). Finally, the hardening law of the slip system  $\alpha$  is determined by Eq. 6.

$$h_{\alpha\beta} = \tau_0 + \int_0^t q_{\alpha\beta} \left[ h_0 \left( 1 - \frac{\tau_c^\beta}{\tau_c^s} \right)^a \right] |\dot{\gamma}^\beta| dt \quad (6)$$

$\tau_0$  is the initial critical shear stress. Considering the quasi-static deformation, the calibrated crystal plasticity parameters of Eq. 6 are  $\tau_0$ ,  $h_0$ ,  $\tau_c^s$  and  $a$ . As the strain rate effect is not considered in this study, the strain rate sensitive parameters,  $\dot{\gamma}_0$  and  $m$ , have been taken from the literature [10].

### 3.3 Dislocation-based crystal plasticity model

In the dislocation-based crystal plasticity model, the dislocation densities ( $\rho_e$ , edge dislocation density and  $\rho_d$ , dipole dislocation density) and their evolution are involved. The Orowan equation is used to define the shear rate of the slip system  $\alpha$ .

$$\dot{\gamma}^\alpha = \rho_e b_s v_0 \exp \left[ -\frac{Q_s}{k_B T} \left\{ 1 - \left( \frac{|\tau_{eff}^\alpha|}{\tau_{sol}} \right)^p \right\}^q \right] \text{sign}(\tau^\alpha) \quad (7)$$

In Eq. 7,  $b_s$ , Burger's vector length for slip, represents the distance and direction of a dislocation in the crystal plane. Dislocation glide velocity,  $v_0$ , represents the velocity of the dislocation, which is dependent on the applied shear stress and temperature.  $Q_s$ , the activation energy for dislocation slip, means the minimum amount of energy needed for a dislocation slip to occur. Previously mentioned

parameters are defined as material constants.  $\tau_{\text{eff}}^{\alpha}$ , refers to the effective resolved shear stress on the  $\alpha$  slip system.  $\tau_{\text{sol}}$ , solid solution strength,  $p$  and  $q$  are fitting parameters.  $k_B$  for Boltzmann constant and  $T$  is temperature.

$\tau_{\text{eff}}^{\alpha}$ , can be defined as:

$$\tau_{\text{eff}}^{\alpha} = \begin{cases} |\tau^{\alpha}| - \tau_{\text{pass}}^{\alpha}, & \text{for } |\tau^{\alpha}| > \tau_{\text{pass}}^{\alpha} \\ 0, & \text{for } |\tau^{\alpha}| < \tau_{\text{pass}}^{\alpha} \end{cases} \quad (8)$$

for Eq. 8,  $\tau^{\alpha}$  is the resolved shear stress.  $\tau_{\text{pass}}^{\alpha}$  is calculated with the following equation (see Eq. 9):

$$\tau_{\text{pass}}^{\alpha} = Gb \left( \sum_{\alpha'=1}^{N_s} \xi_{\alpha\alpha'} (\rho_e^{\alpha'} + \rho_d^{\alpha'}) \right)^{\frac{1}{2}} \quad (9)$$

where,  $G$  and  $\xi_{\alpha\alpha'}$  are material constants, defined as the shear modulus and interaction matrix between  $\alpha$  and  $\alpha'$  slip planes, respectively.

The evolution of edge and dipole dislocation densities on  $\alpha'$  slip system is defined as,

$$\dot{\rho}_e^{\alpha} = \frac{|\dot{\gamma}^{\alpha}|}{b_s \Lambda_s^{\alpha}} - \frac{2\hat{d}^{\alpha}}{b_s} \rho_e^{\alpha} |\dot{\gamma}^{\alpha}| - \frac{2d^{\alpha}}{b_s} \rho_e^{\alpha} |\dot{\gamma}^{\alpha}| \quad (10)$$

$$\dot{\rho}_d^{\alpha} = \frac{2\hat{d}^{\alpha}}{b_s} \rho_e^{\alpha} |\dot{\gamma}^{\alpha}| - \frac{2d^{\alpha}}{b_s} \rho_d^{\alpha} |\dot{\gamma}^{\alpha}| - \rho_d^{\alpha} \frac{4v_{\text{climb}}}{(\hat{d}^{\alpha} - d^{\alpha})} \quad (11)$$

where, the mean free path for slip is defined as  $\Lambda_s^{\alpha}$  and climb dislocation velocity given by,

$$v_{\text{climb}} = \frac{3GD_0\Omega}{2\pi k_B T} \frac{1}{(\hat{d}^{\alpha} + d^{\alpha})} \exp \left( -\frac{Q_c}{k_B T} \right) \quad (12)$$

where, climbing activation volume  $\Omega$  is the fitting parameter to control the volume required by the material in the process of surpassing the energy barrier. The activation energy for climb, the FCC aluminum self-diffusivity coefficients ( $Q_c$ , and  $D_0$ ) are physical constants that define the motion of a single molecule in the same molecule layer.  $\hat{d}^{\alpha}$ , is the maximum slip plane distance required for two dislocations to form a dipole and  $d^{\alpha}$  is the minimum distance required for the elimination of two edge dislocations, which is a fitting parameter.  $\hat{d}^{\alpha}$  is defined in the following equation:

$$\hat{d}^{\alpha} = \frac{3Gb_s}{16\pi|\tau^{\alpha}|} \quad (13)$$

The dislocation mean free path will describe hardening during the straining process. Considering the slip as the deformation mechanism, it can be described as:

$$\frac{1}{\Lambda_s^{\alpha}} = \frac{1}{d} + \frac{1}{\lambda_{\text{slip}}^{\alpha}} \quad (14)$$

$$\frac{1}{\lambda_{\text{slip}}^{\alpha}} = \frac{1}{i_{\text{slip}}} \left( \sum_{\alpha'=1}^{N_s} \xi_{\alpha\alpha'} (\rho_e^{\alpha'} + \rho_d^{\alpha'}) \right)^{\frac{1}{2}} \quad (15)$$

From here two new parameters are obtained.  $d$  is the average grain size of material, and  $i_{\text{slip}}$  is another fitting parameter.

#### 4 Parameter Calibration

In both phenomenological and dislocation-based models, Hooke's law is used for explaining the elastic behavior and the same elastic parameters from the literature [13] have been used. In addition, these two models also share some same settings for the calculation of plastic deformation. For example, the solid solution strength,  $\tau_{\text{sol}}$  or  $\tau_0$ . For the rest of the fitting parameters, Fig. 3 shows a schematic drawing of the calibration process. For both models, the parameters controlling the yield point shall be calibrated first. After the experimental yield strength has been picked in the simulation, the rest hardening parameters can be calibrated according to the flow curve obtained from RD uniaxial tension. Table 1 lists the parameter set used for the phenomenological model.

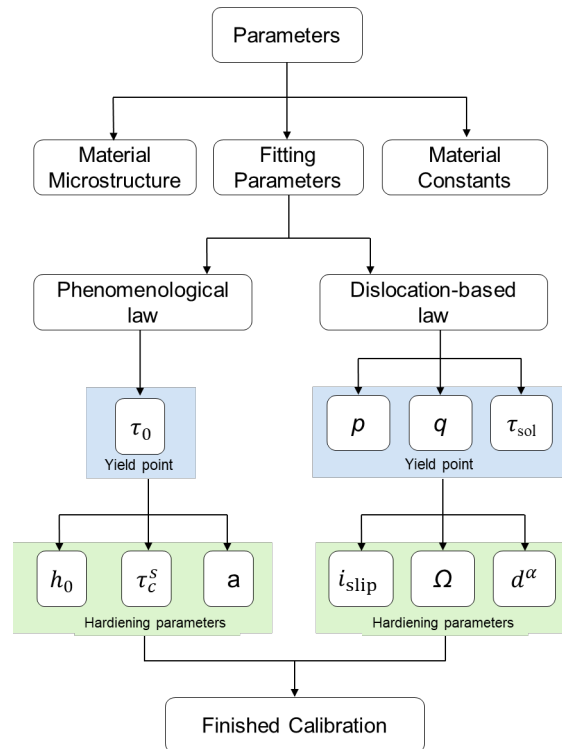


Fig. 3 Followed procedure for parameter calibration.

Table 1 Calibrated parameters for phenomenological CP model

Parameter	$c_{11}$	$c_{12}$	$c_{44}$	$\tau_0$	$h_0$	$\tau_c^s$	$a$
Value	206	118	54	50	600	140	1.6
Unit	GPa	GPa	GPa	MPa	MPa	MPa	-
Parameter	$m$ [13]	$\dot{\gamma}_0$ [13]	$a_{slip}$				
Value	0.05	0.001	1, 1, 1.4, 1.4, 1.4, 1.4				
Unit	-	$s^{-1}$	-				

In addition to the fitting parameters, there are other microstructural parameters and material constants in the dislocation-based model that need to be fixed first. The microstructural parameters, like grain size ( $d$ ) and dipole dislocation density ( $\rho_d$ ), can be analyzed based on the EBSD measurement or obtained from the literature. The rest of the material constants such as Burger's vector, dislocation glide velocity, activation energies for slip and climb, and self-diffusion coefficient are obtained from the literature. All parameters and constants used for the dislocation-based model are summarized in Table 2.

To calibrate all six parameters in the dislocation-based mode, a parametric study has been carried out to investigate the effects of each parameter. The flow curve from tensile testing along RD is used for the parameter calibration. It is indicated that  $p$  influences the yielding point of the material and while the parameter value increases, the yield strength also increases. While the parameter  $q$  has opposite effects with the  $p$  parameter, that is, with the  $q$  value decreasing, the yield strength increases.  $\tau_{sol}$  is the third parameter controlling the yield strength, a higher value for solid solution strength is the equivalence of a higher-yielding point. The  $i_{slip}$  parameter controls the plastic hardening region and leads to the different hardening coefficients. With a lower value of  $i_{slip}$ , the strength of the material will be higher. The effects of  $\Omega$  and  $d^\alpha$  are similar and just control the end part of the flow curve. It can be seen that the calibration process of the dislocation-based model is more difficult than that of the phenomenological model in this section. It involves more parameters and a lot of effort. Therefore a quantitative evaluation criterion is necessary to assess the parameter calibration performance.

Table 2 Material parameters and constants for the dislocation-based CP model

	Description	Value	Unit
Material microstructure parameters	Grain Size ( $d$ )	15.25	$\mu\text{m}$
	Edge dislocation density ( $\rho_e$ ) [9]	1e12	$\text{m}^{-2}$
	Dipole dislocation density ( $\rho_d$ ) [14]	1	$\text{m}^{-2}$
	Burgers vector for slip [9] ( $b_s$ )	2.546e-10	m
	Activation energy for slip [9] ( $Q_s$ )	8.36e-20	J
	Activation energy for climb [9] ( $Q_c$ )	1.876e-19	J
Material Constants	Dislocation glide velocity [14] ( $v_0$ )	1e-4 *	m/s
	Self-diffusion coefficient for Al [15] ( $D_0$ )	6.23e-4	$\text{m}^2/\text{s}$
	C11 [16]	206	GPa
	C12 [16]	118	GPa
	C44 [16]	54	GPa
	Interaction Coefficients [6] ( $\xi\alpha\alpha'$ )	0.122 0.122 0.625 0.07 0.137 0.122	-
Calibrated parameters	Top of the obstacle profile ( $p$ )	0.95	-
	Bottom of the obstacle profile ( $q$ )	1.2	-
	Solid solution strength ( $\tau_{\text{sol}}$ )	50	MPa
	Average dislocation spacing during dislocation travels ( $i_{\text{slip}}$ )	28	-
	Controlling the atomic volume ( $\Omega$ )	4	-
	Controlling the minimum dipole distance ( $d^\alpha$ )	7	-

## 5 Results and Discussion

### 5.1 Prediction on linear loading case

The experimental result in Fig. 4 shows an inhomogeneous plastic deformation, which results in a serrated stress–strain curve. This phenomenon is attributed to the Portevin–Le Chatelier (PLC) effect. This effect, called also dynamic strain aging, is a plastic instability observed in many alloys (such as Al-Mg alloys) when deformed at certain ranges of strain rates and temperatures [17]. The origin of the PLC effect lies in dynamic interactions between mobile solid solution atoms and dislocations.

For analyzing the anisotropic behavior of investigated material, the stress-strain responses of different loading angles have been used for testing the prediction of the two models and comparing their performance (See Fig. 4, as similar results have been obtained in all loading angles, here only 0°/45°/90° results are presented.). It can be seen in the case of the dislocation-based model, early strain rates seem quite good under different loading angles. However, the prediction results of this method are not very good at medium and large strain levels, which may be related to the low quality of parameter calibration in this range. While the phenomenological model shows good adaptability with the loading angle change. It makes a good prediction of large strain values at almost all loading angles. But for the early strain rates, the dislocation-based model can match the experimental data better than the phenomenological model. Particularly for 90°, the stress prediction results at large plastic strain are notably lower than the experimental data. For 45°, both of the phenomenological and dislocation-based models show a good performance for the prediction.

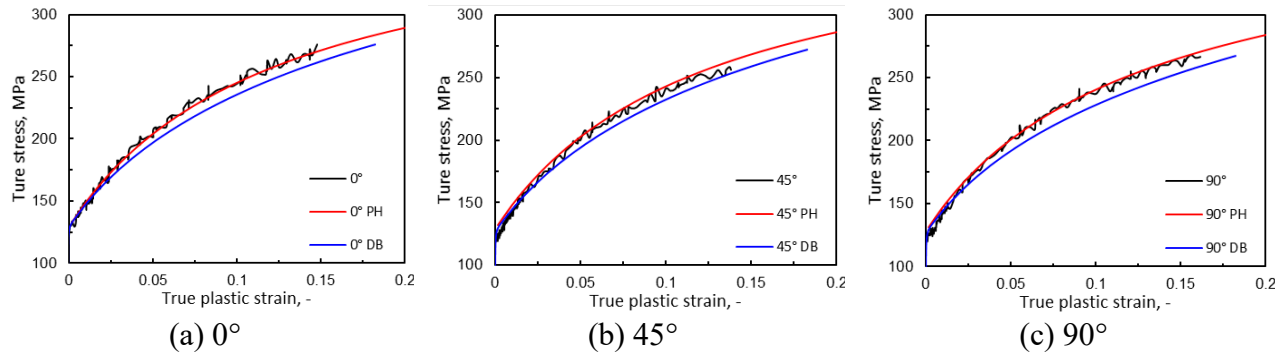


Fig. 4 Anisotropic prediction of phenomenological and dislocation-based CP models.

## 5.2 Prediction on non-linear loading case

In this part, the experimental values were compared with the predicted values of the CP model, as shown in Fig. 6. The experimental data come from two parts, one is the values of the linear tensile test mentioned in the previous section, the other is the values obtained by the tensile test of the SDB that were cut from the material after three kinds of pre-strain of the material. There are three kinds for pre-strain the biaxial, plane strain, and uniaxial pre-strain. The simulation results of the CP model also come from two parts, one is the phenomenological model and the other is the dislocation-based model. In addition, all experiments and simulations have been carried out from both RD and TD directions.

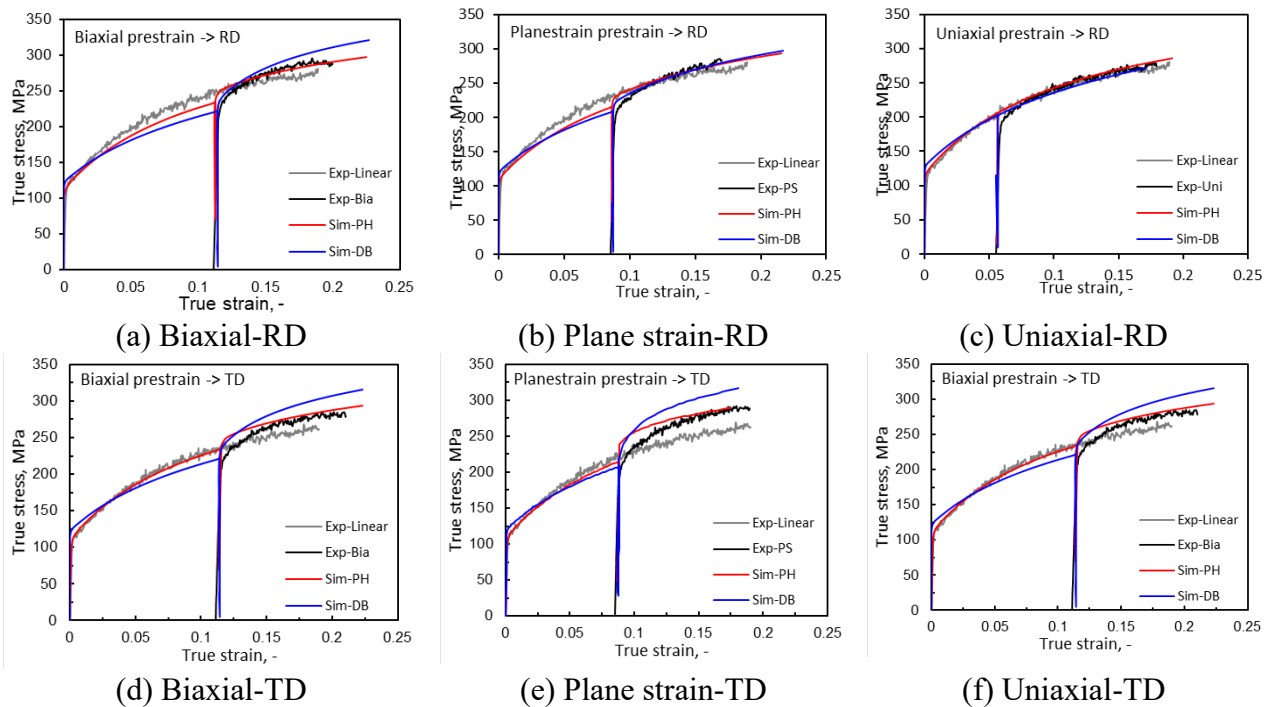


Fig. 5 Non-linear strain paths of phenomenological and dislocation-based CP model.

From the experimental results, it is clear that the stress response has changed after a certain pre-strain loading. Particularly for TD, the stress values of the materials after the three kinds of pre-strain are much higher than the experimental values of linear loading at the same strain level. However, the stress values of the uniaxial strain path are only slightly higher than the initial specimen, especially for the RD uniaxial mini-SDB specimen (Fig. 5 c), as the stress states are identical in two stages.

Overall, both models performed slightly better along the RD. This is because the calibration of parameters has been carried out based on the initial RD linear loading result. For the biaxial loading path (Fig. 5 a), the simulation results of the dislocation-based model have better prediction than the phenomenological model at the smaller strain values. However, it shows a large deviation under large strain conditions. On the contrary, phenomenology shows higher quality at larger strain values. For



the uniaxial loading path (Fig. 5 c), there is a little difference between the experiment values under uniaxial pre-strain loading and the initial tensile test values. Both of the phenomenological and dislocation-based models did a good prediction under this loading condition.

Both models show large deviations in the prediction along the TD. However, the phenomenological model performs well in the prediction of yield point and the late deformation stage, while the shape and extent tendency of the dislocation-based model prediction is very similar to those of the experimental value curves, indicating that the dislocation-based model may have a better adaptability in the prediction of the second stage loading. Hence, it is possible to expect a better prediction accuracy of the dislocation-based model if an improved parameter calibration could be conducted.

Currently achieved results indicate that the dislocation-based model is highly sensitive to the change of the stress states. With stress state changes, the slip system would change, and it has a significant influence on stress-strain responses. In this study, the parameter of interaction between slip systems in the dislocation-based model is referred to Kubin, et al. [18], which is not totally the same as the interaction parameter used in the phenomenological model. There are still other examples of slip interaction parameters that can be verified in the future [14]. For the comparison between two CP models, the effect of the slip interactive parameter shall be considered. Generally, the non-linear strain path prediction has shown bigger errors than the linear strain path, but it has been able to capture the general tendencies in the material behavior. Some further research could be carried out to refine the predictions. For example, the quality of parameter calibration shall be improved, and multiple standards such as R-value and hardening curve would be involved to guarantee the accuracy of the model.

## 6 Remarks

In this study, the plastic deformation behavior of a single-phase aluminum alloy under linear and non-linear loading has been investigated by both experimental and numerical methods. Two crystal plasticity models, phenomenological model, and dislocation-density-based model respectively, have been employed, and their performance has been compared in terms of the parameter calibration procedure and the prediction results. Several remarks can be drawn in the following.

- The dislocation-density-based model requires an extremely complex parameter calibration process. This model includes more microstructure information compared with the phenomenological one and it has more parameters affecting different properties of the material, such as grain size and dislocation density. It needs more effort to calibrate the dislocation-based model. A quantitative evaluation criterion is necessary to assess the parameter calibration performance.
- Both models have good performance for the prediction of linear loading cases. It is concluded that the dislocation-based model shows a better yield strength prediction while the phenomenological performed with higher quality at large strain levels, which is consistent with their parameter calibration performance. Therefore, the identical parameter calibration performance should be guaranteed first to further compare the model performance.
- The non-linear strain path prediction has shown bigger errors, but it has been able to capture the general tendencies in the material behavior. Some further research could be carried out to refine the predictions. For example, the quality of parameter correction will be improved, and studying the influence of slip system under different stress conditions.
- To evaluate the model performance, other mechanical and material responses, such as R-value and texture evolution should also be focused on in the future. hardening curve.

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**References**

- [1] L.Z. W.S Miller, J Bottema, A.J Wittebrood, P De Smet, A Haszler, A Vieregge, *Materials Science and Engineering: A* 280 (2000) 37-49.
- [2] S. Hou, Y.L. Zhu, Y.L. Wang, H.X. Sun, *Rare Metal Materials and Engineering* 46 (2017) 3760-3766.
- [3] T. Hasebe, M. Sugiyama, H. Adachi, S. Fukutani, M. Iida, *Materials Transactions* 55 (2014) 779-787.
- [4] T. Skippon, C. Mareau, M.R. Daymond, *Journal of Applied Crystallography* 45 (2012) 627-643.
- [5] A.B. Lopes, F. Barlat, J.J. Gracio, J.F.F. Duarte, E.F. Rauch, *International Journal of Plasticity* 19 (2003) 1-22.
- [6] J.W. Yoon, F. Barlat, J.J. Gracio, E. Rauch, *International Journal of Plasticity* 21 (2005) 2426-2447.
- [7] J. Wang, T.J. Chen, C.Z. Zhou, *Tribology International* 156 (2021).
- [8] W.L. Junhe Lian, Xabier Gastañares, Rongfei Juan, Joseba Mendiguren, *International Journal of Multiphase Flow* Under review (2021).
- [9] M.S. Pham, M. Iadicola, A. Creuziger, L. Hu, A.D. Rollett, *International Journal of Plasticity* 75 (2015) 226-243.
- [10] W. Liu, J. Lian, N. Aravas, S. Münstermann, *International Journal of Plasticity* 126 (2020) 102614.
- [11] F. Bachmann, R. Hielscher, H. Schaeben, *Texture and Anisotropy of Polycrystals Iii* 160 (2010) 63-+.
- [12] F. Roters, P. Eisenlohr, L. Hantcherli, D.D. Tjahjanto, T.R. Bieler, D. Raabe, *Acta Materialia* 58 (2010) 1152-1211.
- [13] S.M. Rassoulinejad-Mousavi, Y.J. Mao, Y.W. Zhang, *Journal of Applied Physics* 119 (2016).
- [14] S.L. Wong, M. Madivala, U. Prahl, F. Roters, D. Raabe, *Acta Materialia* 118 (2016) 140-151.
- [15] E.D. Cyr, A. Brahme, M. Mohammadi, R.K. Mishra, K. Inal, *Materials Science and Engineering a-Structural Materials Properties Microstructure and Processing* 727 (2018) 11-28.
- [16] M.D. Pandey, W.-C. Xie, L. Xu, *Advances in engineering structures, mechanics & construction: proceedings of an International Conference on Advances in Engineering Structures, Mechanics & Construction, held in Waterloo, Ontario, Canada, May 14-17, 2006*, Springer, Dordrecht, Netherlands, 2006, p. xvi, 866 p.
- [17] A. Chatterjee, P. Mukherjee, N. Gayathri, P. Barat, A. Barat, A. Sarkar, *Bulletin of Materials Science* 34 (2011) 1113-1117.
- [18] L. Kubin, B. Devincre, T. Hoc, *Acta Materialia* 56 (2008) 6040-6049.