

Models for Impurity Incorporation during Vapor-Phase Epitaxy

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Abstract. Impurity incorporation during vapor-phase epitaxy on stepped surfaces was modeled by classifying rate-limiting processes into i) surface diffusion, ii) step kinetics, and iii) segregation. Examples were shown for i) desorption-limited Al incorporation during chemical vapor deposition (CVD) of (0001) SiC, ii) preferential desorption of C atoms from kinks during CVD of Al-doped (000-1) SiC, and iii) segregation-limited C incorporation during metalorganic vapor-phase epitaxy of (0001), (000-1), and (10-10) GaN.

Introduction

Impurity incorporation during vapor-phase epitaxy has been modeled via, for example, site competition [1,2] and surface vacancies [3,4]. The latter, however, cannot explain the variation in impurity doping around facets [5]. Moreover, in the cases of homoepitaxial growths of SiC and GaN, misoriented substrates are often used for polytype [6] and doping-uniformity [7] controls, respectively. Accordingly, we modeled impurity incorporation during step-flow growth by taking Al-doped SiC and C-doped GaN, as examples. We believe the models should be beneficial for determining allowable off-angle variations for desired doping-level uniformities in advanced devices. Although Al was chosen due to the availability of a thermodynamic model [8], N doping for SiC could be similarly treated under the assumption of the N segregation coefficient being unity [9].

Proposed Models

Impurity incorporation during vapor-phase epitaxy on stepped surfaces was modeled by classifying rate-limiting processes into i) surface diffusion [10], ii) step kinetics [11], and iii) segregation [12] (Table I).

i) **Desorption** limits impurity incorporation at step-edges when surface diffusion length λ is less than a half of the average inter-step distance, λ_0 . This should be the case with incorporation of Al, whose λ was estimated to be less than 2 nm at 1550°C [10], into stepped 4H-SiC (0001). This is due to relatively large λ_0 (eg., 7.2 nm for $\theta = 8^\circ$) originating from four-bilayer-high steps [13]. Based on the Burton–Cabrera–Frank (BCF) theory [14], we derived the following equation for x in $\text{Al}_x\text{Si}_{1-x}\text{C}$ [10]:

$$F_{\text{Al}} / x = [\gamma P_{\text{Si}}^e / K (2 \pi m_{\text{Al}} k_B T_g)^{1/2}] + [F_{\text{Si}} - P_{\text{Si}}^e / (2 \pi m_{\text{Si}} k_B T_g)^{1/2}] [\lambda_0 / 2\lambda_{\text{Al}} \tanh (\lambda_0 / 2\lambda_{\text{Al}})], \quad (1)$$

where F_i , P_i^e , and m_i ($i = \text{Al}, \text{Si}$) are, respectively, the incident flux, equilibrium vapor pressure, and mass of i atom, K and γ are, respectively, the equilibrium constant and activity coefficient for AlC, T_g is growth temperature, and k_B is Boltzmann's constant. Eq. (1) explains why x was independent of the off-angle θ (ranging from 2° to 8°) when the C/Si ratio, r , was small (i.e., 1.8 [15]); due to large P_{Si}^e , the first term in the right-hand side, which corresponds to the Al desorption flux, became

dominant (solid line in Fig. 1). Eq. (1) also explains why x increased with θ when r was large (i.e., 4–6 [16]); due to small P_{Si}^e , the second term in the right-hand side, which corresponds to the Al flux incorporated into the solid, became so large that x increased with the step density on the surface (dashed and dotted lines in Fig. 1).

Table I. Rate-limiting processes of impurity incorporation during vapor-phase epitaxy.

Classification of host-atom desorption from kinks	Surface-diffusion length of impurity atoms	
	Less than $\lambda_o/2$	Much larger than $\lambda_o/2$
Preferential desorption of host atoms from kinks	Surface diffusion	Step kinetics
Negligible desorption of host atoms from kinks		Segregation

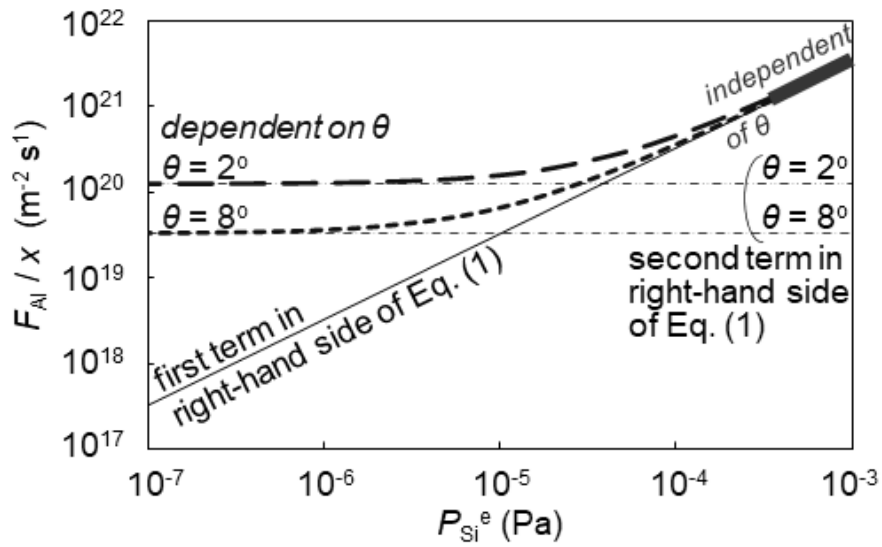


Fig. 1. F_{Al}/x , calculated as first term (solid line) and second term (dashed and dotted lines) in right-hand side of Eq. (1), as a function of equilibrium vapor pressure of Si, with assumptions of T_g of 1550°C, growth rate of 1.3 $\mu\text{m/h}$, and λ_{Al} of 2.0 nm.

ii) Preferential desorption of host atoms from kinks limits impurity incorporation at kinks even when $\lambda \gg \lambda_o/2$. This should be the case with incorporation of Al into 4H-SiC (000-1) that has one-bilayer-high steps [13]. We assume that a C atom making two bonds with a Si atom stays at kinks, while that a C atom making one bond with a Si atom easily desorbs from kinks [Fig. 2(a)].

Since r is typically small (eg., $r \leq 6$ [16]), some surface-diffusing Al atoms that arrive at kinks keep waiting (for an average time τ_c) until C atoms make one bond with Si atoms at kinks [Fig. 2(b)] before they are incorporated into the solid [Fig. 2(c)]. Based on the reported experimental results [16], surface Al concentration n_{Al} (normalized by the mean residence time τ_{Al}) was calculated (Fig. 3). n_{Al} in the vicinity of step-edges (i.e., local minima in Fig. 3) on (000-1) is much larger than that on (0001), indicating longer τ_c on (000-1).

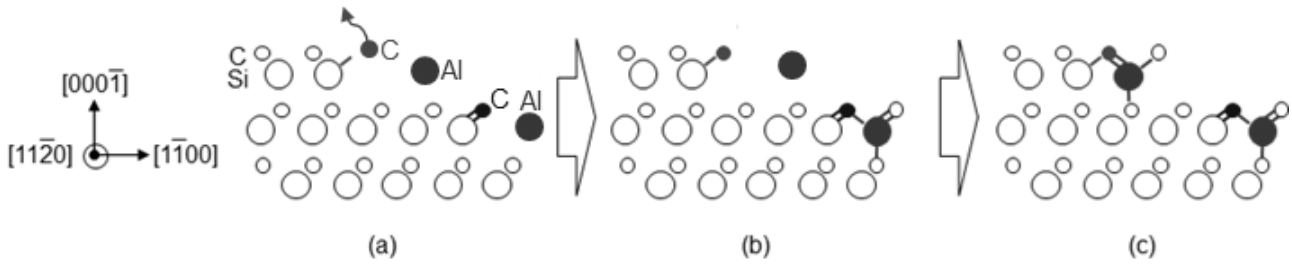


Fig. 2. Schematic illustrations of (a) preferential desorption of a C atom having one bond with a Si atom, (b) adsorption of a C atom to a dangling bond of a Si atom and bonding of an Al atom to two C atoms, and (c) bonding of another Al atom to three C atoms at kinks on 4H-SiC (000-1).

iii) **Segregation** limits impurity incorporation even when $\lambda \gg \lambda_o/2$ and desorption of host atoms from kinks is negligible. This should be the case with incorporation of C into GaN that is typically grown with the N/Ga ratio exceeding 1000 [17–19]; namely, soon after a N atom making one bond with a Ga atom desorbs from kinks, another N atom makes one bond with the Ga atom. When the length of time before the C concentration at the step-edge site reaches its equilibrium value, τ_{step} , is much smaller than the meantime until a C atom incorporated at kinks moves through the step-edge site to the surface site, τ , the C concentration in the solid can be expressed as [20]

$$N = N_{\text{surf}} + (N_{\text{step}} - N_{\text{surf}}) \exp(-D / V_{\text{step}} a), \quad (2)$$

where N_{surf} and N_{step} are, respectively, the equilibrium C concentrations at the surface site and at the step-edge site, D is the diffusion coefficient in the solid, V_{step} is the average step velocity, and a is the lattice constant. As shown in Fig. 4, the results for (0001) [17], (000-1) [18], and (10-10) [19] growths are well reproduced with D of $2 \times 10^{-13} \text{ cm}^2/\text{s}$ that agrees with the experimentally determined value [21].

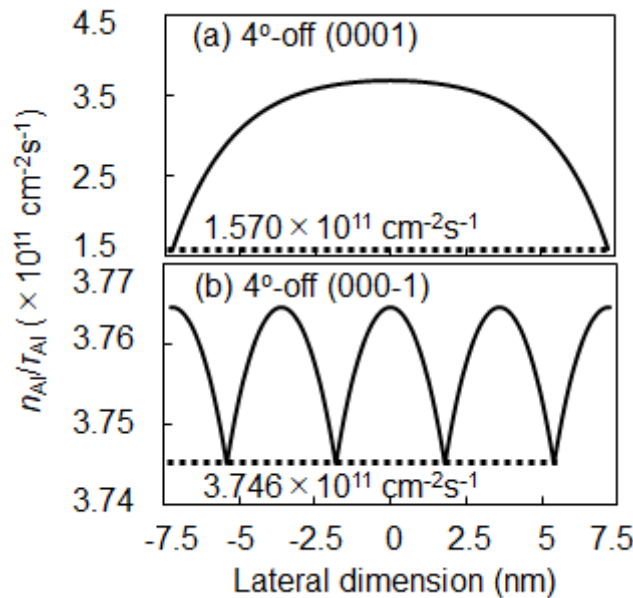


Fig. 3. Distribution of $n_{\text{Al}}/\tau_{\text{Al}}$ calculated with assumptions of T_g of 1550°C, growth rate of 1.3 $\mu\text{m}/\text{h}$, r of 6, and λ_{Al} of 2.0 nm.

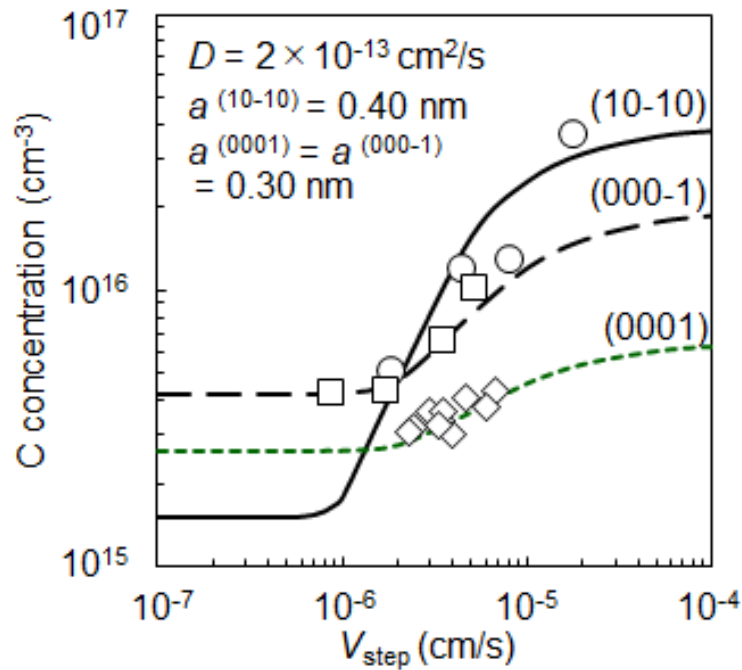


Fig. 4. Step-velocity dependences of carbon concentrations fitted to the reported results [17–19].

Summary

Impurity incorporation during step-flow growth was modeled and exemplified by SiC: Al and GaN: C cases. We believe the proposed models should contribute to determining allowable off-angle variations for desired doping-level uniformities in advanced SiC and GaN devices.

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References

- [1] D.J. Larkin, P.G. Neudeck, J.A. Powell and L. G. Matus, Appl. Phys. Lett. 65 (1994) 1659–1661.
- [2] R. Wang and A.J. Steckl, J. Cryst. Growth 312 (2010) 680–684.
- [3] G. Ferro and D. Chaussende, Sci. Rep. 7 (2017) 43069-1–43069-11.
- [4] G. Ferro, D. Chaussende and N. Tsavdaris, J. Cryst. Growth 507 (2019) 338–343.
- [5] K. Yokomoto, M. Yabu, T. Hashiguchi and N. Ohtani, J. Appl. Phys. 128 (2020) 135701-1–135701-9.
- [6] N. Kuroda, K. Shibahara, W. S. Yoo, S. Nishino and H. Matsunami, Ext. Abstr. 19th Conf. Solid State Devices & Materials, 1987, 227–229.
- [7] F. Horikiri, Y. Narita, T. Yoshida, T. Kitamura, H. Ohta, T. Nakamura and T. Mishima, IEEE Trans. Semicond. Manufact. 30 (2017) 486–492.
- [8] A.S. Segal, S.Yu. Karpov, A.V. Lobanova, E.V. Yakovlev, K. Hara and M. Naito, Mater. Sci. Forum 821–823 (2015) 145–148.
- [9] K. Mochizuki and T. Mishima, Jpn. J. Appl. Phys. 60 (2021) 018001-1–018001-3.
- [10] K. Mochizuki and T. Mishima, Jpn. J. Appl. Phys. 59 (2020) 088003-1–088003-4.

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- [11] K. Mochizuki, F. Horikiri, H. Ohta and T. Mishima, *Jpn. J. Appl. Phys* 59 (2020) 068001-1–068001-3.
 - [12] K. Mochizuki, F. Horikiri, H. Ohta and T. Mishima, *Jpn. J. Appl. Phys* 60 (2021) 018002-1–018002-3.
 - [13] T. Kimoto, A. Itoh, H. Matsunami and T. Okano, *J. Appl. Phys.* 81 (1997) 3494–3500.
 - [14] W.K. Burton, N. Cabrera, and F.C. Frank, *Philos. Trans. R. Soc. London, Ser. A* 243 (1951) 299-358.
 - [15] Information on <https://tel.archives-ouvertes.fr/tel-01466713/document>
 - [16] T. Yamamoto, T. Kimoto and H. Matsunami, *Mater. Sci. Forum* 264–268 (1998) 111–114.
 - [17] F. Horikiri, Y. Narita, T. Yoshida, T. Kitamura, H. Ohta, T. Nakamura and T. Mishima, *Jpn. J. Appl. Phys.* 56 (2017) 061001-1–061001-6.
 - [18] K. Nagamatsu, Y. Ando, T. Kono, H. Cheong, S. Nitta, Y. Honda, M. Pristovsek and H. Amano, *J. Cryst. Growth* 512 (2019) 78–83.
 - [19] H. Yamada, H. Chonan, T. Takahashi and M. Shimizu, *Jpn. J. Appl. Phys.* 57 (2018) 04FG01-1–04FG01-5.
 - [20] T. Nishinaga, C. Sasaoka and K. Park, *Jpn. J. Appl. Phys* 28 (1989) 836–840.
 - [21] X.A. Cao, R.G. Wilson, J.C. Zolper, S.J. Pearton, J. Han, R.J. Shul, D.J. Rieger, R.K. Singh, M. Fu, V. Scarvepalli, J.A. Sekhar and J.M. Zavada, *J. Electron. Mater.* 28 (1999) 261-265.